C 1.5 Radial Expansion Wave

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I. Code Description

The solver relies a standard Discontinuous Galerkin discretization to achieve high-orders of accuracy. The spatial fluxes are approximated using the Roe approximate Riemann solver.[1] All integrals are evaluated analytically during a pre-processing stage using a symbolic manipulation package. The fluxes are computed from the conservative variables using a modal representation of the polynomials rather than a traditional nodal method. A more detailed description of the discretization and solver can be found in Refs. 2 and 3. The code is currently capable of up to 4^{th}-order polynomial approximations of the solution, which yields a 5^{th}-order accurate scheme. There is no inherent limitation that prevents polynomial approximations of higher order.

All calculations are in double precision. Parallel capabilities are forthcoming.

II. Case Summary

The solution was initialized to the analytical initial solution provided using $\gamma = 3$, and extrapolation boundary conditions were imposed on all boundaries. Time integration was achieved with a standard 4^{th}-order Runge-Kutta explicit scheme using a time step calculated using CFL number obtained by

$$CFL = \frac{0.25}{2P+1}.$$

Final entropy error values were evaluated at $t = 2$. Only 2D calculations were performed. Entropy error was computed using the definition stated as part of the case.

A summary of the exucation times and the average execution time of TauBench is shown in Table 1. The code was executed on an a machine with two 6 cores Intel i7 Xeon X5660 2.8 GHz processors. Only one core was used for all calculations.

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Run	Execution Time (s)
1	7.815
2	7.674
3	7.698
4	7.962
5	7.669
6	7.712
7	7.691
8	7.953
9	7.676
10	7.668
Avg:	7.752

Table 1: Taubench Results

Timing of 100 right hand side (RHS) evaluations with 250,000 degrees of freedom (DOF) are summarized in Table 2. The RHS was evaluated on a uniform rectangular mesh. The solution was initialized to a free-stream flow with slip walls on the lower and upper surfaces. As the order of the polynomial approximation increased, the cell size of the mesh was reduced in order to maintain constant DOF. The 100 RHS evaluations were achieved by marching with a 4^{th}-order Runge-Kutta scheme 25 iterations.

Polynomial Order	Mesh Size	DOF	Work
0	500x500	250,000	6.24
1	250x250	250,000	2.70
2	167x167	251,001	2.58
3	125x125	250,000	3.32
4	100x100	250,000	4.92

Table 2: RHS Timing

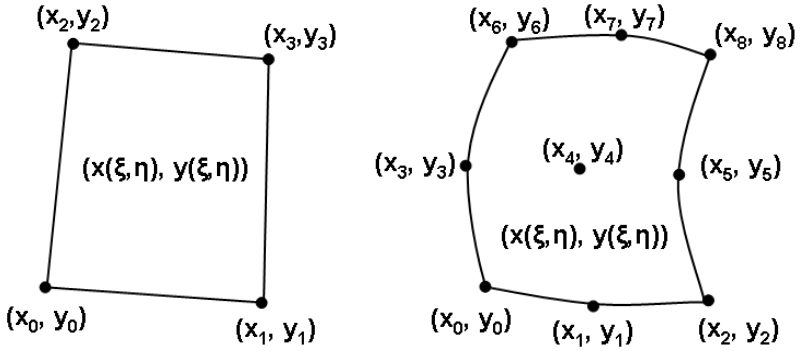
III. Meshes

The DG solver is formulated for structured meshes. These meshes can be generated using any traditional mesh generator. After reading the mesh coordinates, the solver generates cell local polynomial representations, $(x(\xi,\eta), y(\xi,\eta))$, of the cell coordinates. The geometric polynomial mapping of the cell coordinates is formulated as a sum involving the same test functions, ψ, as used in the DG discretization of the governing equations. Hence,

$$\begin{aligned} x(\xi,\eta) &= \sum_{i}^{N_g}\sum_{j}^{N_g} x_{ij}\psi_{ij}(\xi,\eta) \\ y(\xi,\eta) &= \sum_{i}^{N_g}\sum_{j}^{N_g} y_{ij}\psi_{ij}(\xi,\eta) \end{aligned} \tag{1}$$

where x_{ij} and y_{ij} are the coefficients of the expansion, and $\psi_{ij}(\xi,\eta) = P_i(\xi)P_j(\eta)$, where $P_i(\xi)$ are the Legendre polynomials. The coefficients are found by equating the expansion with the associated cell nodal

values. The process is repeated for the y coordinate to obtain the final polynomial mapping of the cell. Additional points are required to establish the polynomial representation as shown for a quadratic cell in Fig. 1b. This implies that the grid must consist of $N_g n+1$ points along a coordinate line to generate n cells with a geometric mapping of order N_g along that line.



(a) Linear cell (b) Quadratic cell

Figure 1: Nodal and modal representation of cells.

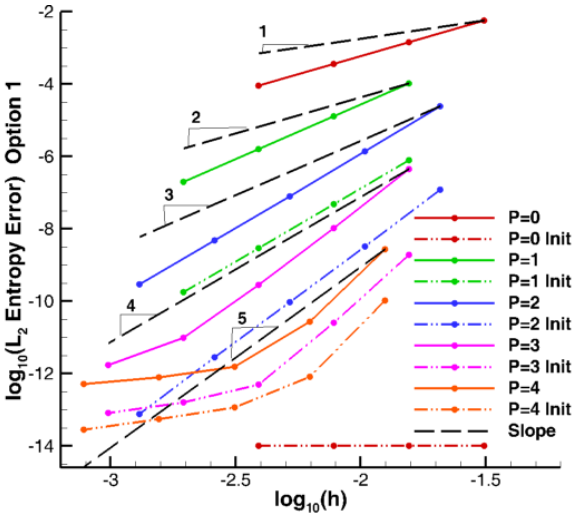
The meshes used uniform cell sizes for the grids on the specified domain of $[-4,4]\times[-4,4]$. The mesh cell sizes used for the calculations are summarized in Table 3. The uniform grids spacing of the mesh generation algorithm automatically achieves a uniform grid refinement.

16x16	32x32	64x64	128x128	256x256

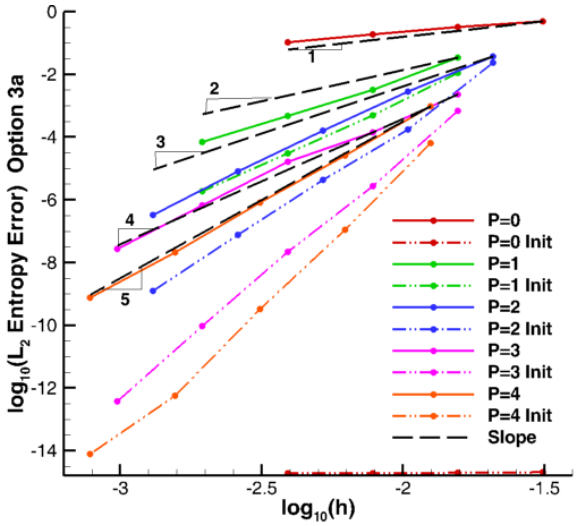
Table 3: Mesh Sizes Used for the Calculations

IV. Results

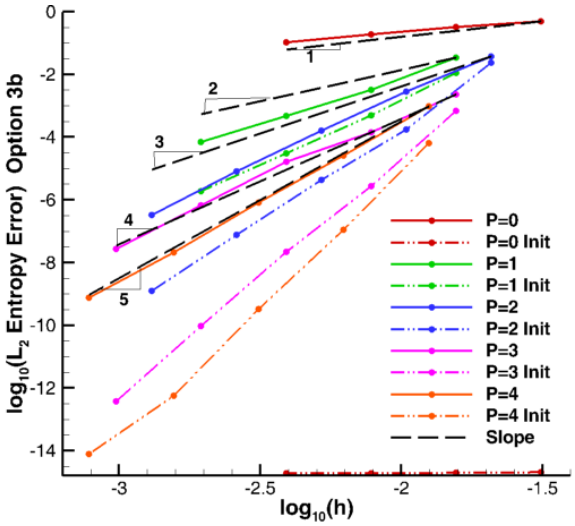
The convergence rate of three types of the L_2-error options are shown in Fig. 2. The convergence rate of the solution achieves the expect $O\left(h^{P+1}\right)$, and is closer to $O\left(h^{p+2}\right)$. L_2-errors vs. Work are shown in Fig. 3. With few exceptions, a higher degree polynomial achieves lower L_2-errors with less work. The L_2-errors history of the solver for the different meshes are shown in Figs. 4, 5, and 6.



(a) Option 1

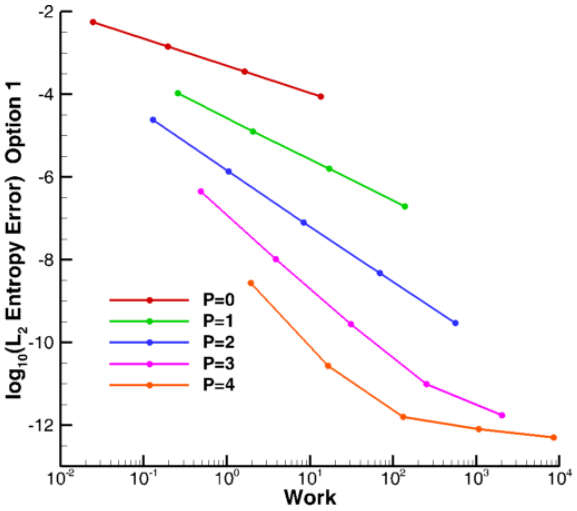


(b) Option 3a

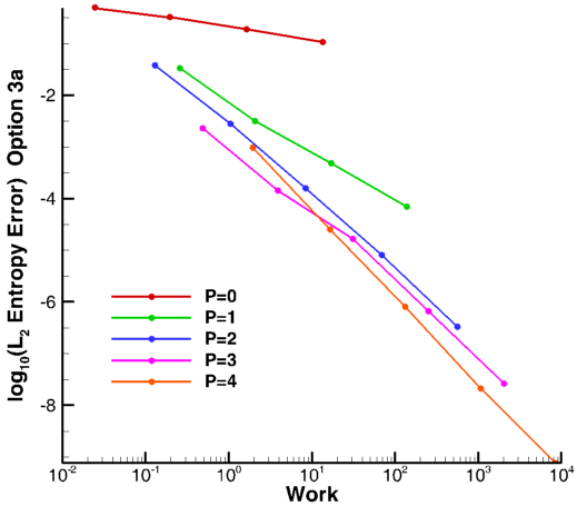


(c) Option 3b

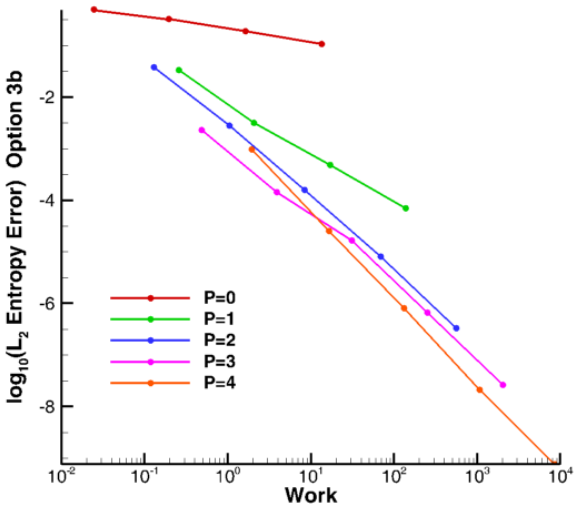
Figure 2: H-Convergence



(a) Option 1

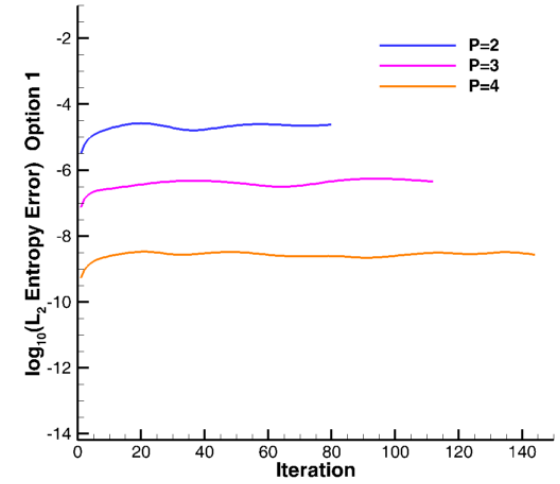


(b) Option 3a

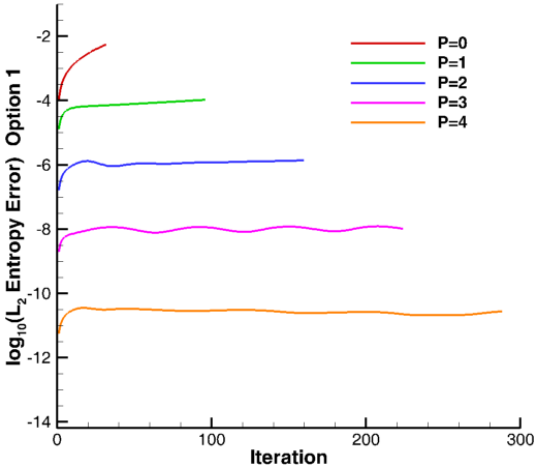


(c) Option 3b

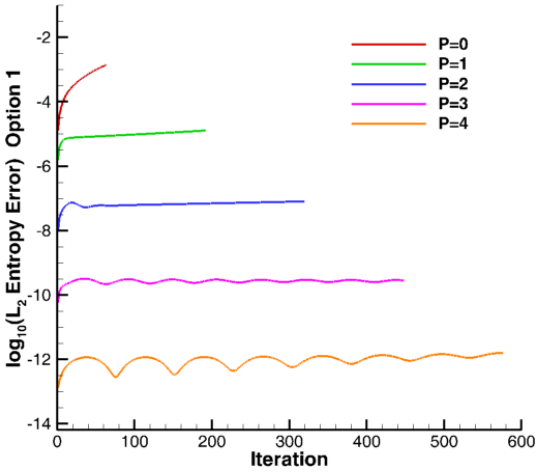
Figure 3: Work



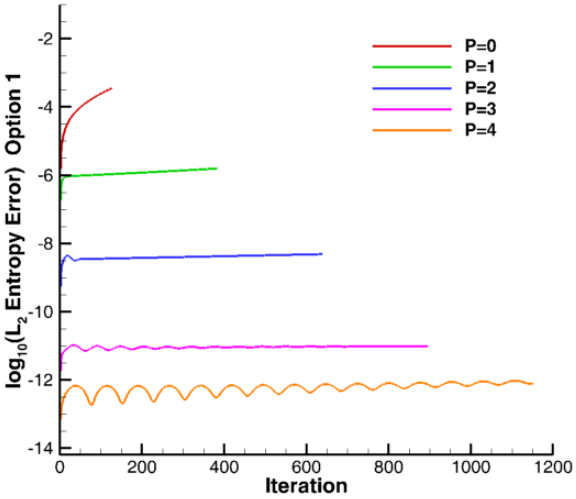
(a) 16x16



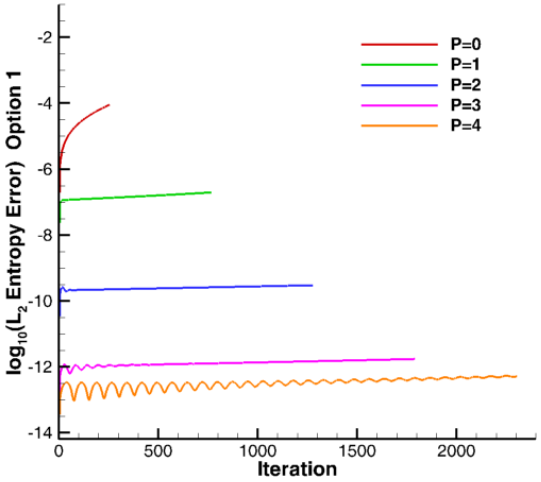
(b) 32x32



(c) 64x64

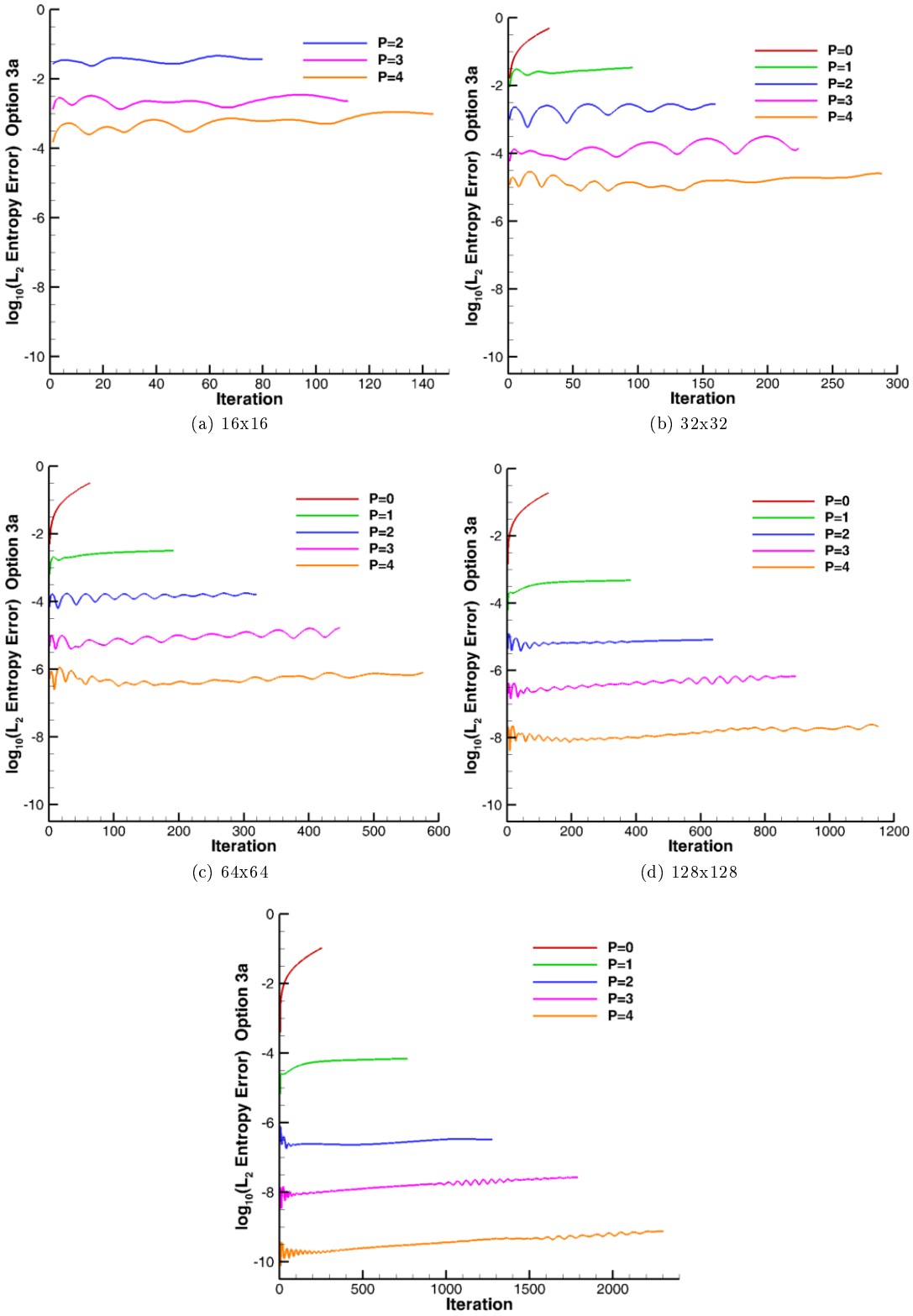


(d) 128x128



(e) 256x256

Figure 4: Iteration History of Option 1 Entropy Error



(a) 16x16

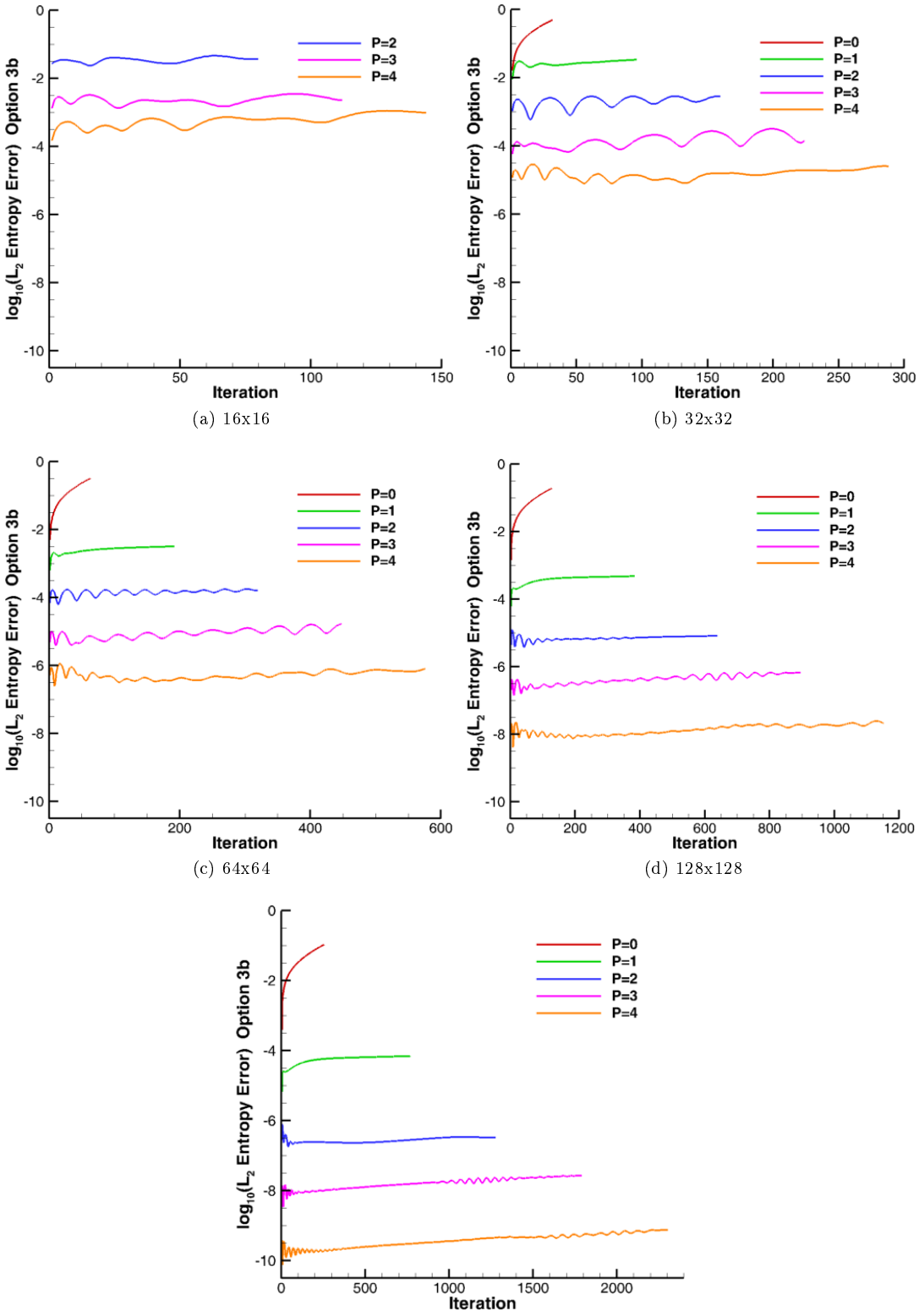
(b) 32x32

(c) 64x64

(d) 128x128

(e) 256x256

Figure 5: Iteration History of Option 3a Entropy Error



(a) 16x16

(b) 32x32

(c) 64x64

(d) 128x128

(e) 256x256

Figure 6: Iteration History of Option 3a Entropy Error

References

[1]Vatsa, V. N. and Wedan, J. L. T. B. W., "Navier-Stokes computations of prolate spheroids at angle of attack," AIAA-Paper 1987-2627, 1987.

[2]Galbraith, M. C., Orkwis, P. D., and Benek, J. A., "Automated Quadrature-free Discontinuous Galerkin Method Applied to Viscous Flows," AIAA-Paper 2011-493, 2011.

[3]Galbraith, M. C., Orkwis, P. D., and Benek, J. A., "Extending the Discontinuous Galerkin Scheme to the Chimera Overset Method," AIAA-Paper 2011-3409, 2011.